Optical Readout TPC (O-TPC) For a Study of Stellar Helium Burning at HIyS

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- 1. The Collaboration
- <u>The Experiment (Briefly):</u> Oxygen Formation in Stellar Helium Burning ¹²C(α,γ)¹⁶O Reaction With Real Photons
- **3.** The Detector (O-TPC)
- 4. Performance of O-TPC
- 5. Future Developments

TUNL, Duke, 12 July 2007

The Laboratory for Nuclear Science At Avery Point



The O-TPC at HIYS Collaboration:

<u>UConn :</u>* M. Gai T.J. Kading L. Weissman A.H. Young P.N. Seo (anticipated) Unnamed Grad Student

Yale:** G.F. Burkhard D.F. Rubin

<u>PTB, Braunschweig:</u> *** B. Bromberger V. Dangendorf K. Tittelmeier

Weizmann, Israel: ** A. Breskin R. Chechik M. Klin

<u>UCL, LLN, Belgium:</u>*** Th. Delbar

- * Supported by US Department of Energy
- ** Supported by the American Committee on Weizman Yale-Weizmann Collaboration
- *** In Kind Contribution, Optical Readout System

TUNL, Duke: M.W. Ahmad S. Stave H.R. Weller

UHartford: J.E. McDonald

<u>GCSU:</u> R.H. France III <u>NGCSU:</u>* R.M. Prior M.C. Spraker



HELIUM BURNING IN (MASSIVE) STARS



III. $\alpha + {}^{16}O = {}^{20}Ne$

(NEGLIGIBLE)



CENTRAL DENSITY (gm/cc)

"Normal" main-sequence companion

Rotation

Mass-transfer stream

White dwarf

(a)

Accretion disk

Alexander Pala Strate Colors a star

"Hot

spot"

Lagrange point

> Roche lobe of companion

Explosion of a White Dwarfs (Defl., Delayed Det. & Merger)



Deflagration: Energy transport by heat conduction over the front, v <<v(sound)=> ignition of unburned fuel (C/O)
Detonation: ignition of unburned fuel by compression, v = v(sound)
Rem1: Pre-expansion depends on the amount of burning. The rate of burning hardly changes the final structure for DD-models (Dominguez et al. ApJ 528, 590)

Rem.2: HeDs (sub-MCh)



- disagree with LCs and spectra (Nugent et al. 96, Hoeflich et al. 96)

Peter Hoeflich (2002)

INFLUENCE ON LIGHT CURVES (0-60 Days)

DD21c: C/O=1/1; Z=0.02 (solar)

DD23c: C/O=2/3; Z=0.02 (solar)

DD24c: C/O=1/1; Z=0.0067 (solar/3)

Bolometric Light Curves



C/O Ratio of the WD

- Maxima \approx 2-3 days later (i.g. 1-5 days)

- Peak to 'Tail' ratio changes by $\approx 0.3^m$ Metalicity Z - negligible

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OFF SET in M(dM15) dM(V) \simeq 0.1 dt(rise)
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The chlorine detector must be maintained in low-level operation until the chlorine and gallium detectors can be operated at full level simultaneously. Otherwise endless conjecture concerning time variations in the solar neutrino flux will ensue. Moreover, the results of the gallium observations may uncover information that has been overlooked in the past chlorine observations.

The CNO cycle operates at the higher temperatures which occur during hydrogen burning in main sequence stars somewhat more massive than the sun. This is the case because the CNO cycle reaction rates rise more rapidly with temperature than do those of the pp chain. The cycle is important because ${}^{13}C$, ${}^{14}N$, ${}^{15}N$, ${}^{17}O$, and ${}^{18}O$ are produced from ${}^{12}C$ and ${}^{16}O$ as seeds. The role of these nuclei as sources of neutrons during helium burning is discussed in Sec. V.

V. THE SYNTHESIS OF ¹²C AND ¹⁶O AND NEUTRON PRODUCTION IN HELIUM BURNING

The human body is 65% oxygen by mass and 18% carbon, with the remainder mostly hydrogen. Oxygen (0.85%) and carbon (0.39%) are the most abundant elements heavier than helium in the sun and similar main se-

quence stars. It is little wonder that the determination of the ratio ${}^{12}C/{}^{16}O$ produced in helium burning is a problem of paramount importance in Nuclear Astrophysics. This ratio depends in a fairly complicated manner on the density, temperature, and duration of helium burning, but it depends directly on the relative rates of the $3\alpha \rightarrow {}^{12}C$ process and the ${}^{12}C(\alpha,\gamma) {}^{16}O$ process. If $3\alpha \rightarrow {}^{12}C$ is much faster than ${}^{12}C(\alpha,\gamma) {}^{16}O$, then no ${}^{16}O$ is produced in helium burning. If the reverse is true, then no ${}^{12}C$ is produced. For the most part the subsequent reaction ${}^{16}O(\alpha,\gamma) {}^{20}Ne$ is slow enough to be neglected.

There is general agreement about the rate of the $3\alpha \rightarrow {}^{12}\text{C}$ process, as reviewed by Barnes (1982). However there is a lively controversy at the present time about the laboratory cross section for ${}^{12}\text{C}(\alpha,\gamma)$ ${}^{16}\text{O}$ and about its theoretical extrapolation to the low energies at which the reaction effectively operates. The situation is depicted in Figs. 4, 5, and 6, taken with some modification from Langanke and Koonin (1983), Dyer and Barnes (1974), and Kettner *et al.* (1982). The Caltech data obtained in the Kellogg Laboratory is shown as the experimental points in Fig. 4, taken from Dyer and Barnes (1974), who compared their results with theoretical calculations by Koonin, Tombrello, and Fox (1974). The Münster data are shown as the experimental points in Fig. 5, taken from

Helium Burning:

 $3\alpha \rightarrow 12C \text{ Known}$ $\alpha + 12C \rightarrow 160 \quad ???$ $\boxed{C/O = ?}$

 $12C(\alpha, \gamma)16O (E_{cm} = 300 \text{ keV})$

 $\sigma(\alpha, \gamma) = S/E \times e^{-2\pi \eta}$ $(\eta = e^{2Z_1Z_2/\hbar \upsilon} = Z_1Z_2\alpha/\beta)$

Astrophysical Cross Section Factor (P and D waves)







Physics Today 55:12(2002)26





¹⁶O +
$$\overrightarrow{\gamma} \rightarrow \alpha$$
 + ¹²C *
E _{γ} = 8.0 - 10.0 MeV (±1%)

 $\begin{aligned} & \left(\sigma(\alpha, \gamma_{cascade}) \right) < 4\% \text{ at } 300 \text{ keV} \\ & \sigma(\gamma, \alpha) = \frac{(2S_1 + 1)(2S_2 + 1)}{2(2S_4 + 1)} \times \frac{k_{\alpha}^2}{k_{\gamma}^2} \times \sigma(\alpha, \gamma) \\ & \sigma(\overrightarrow{\gamma}, \alpha) = \frac{(2S_1 + 1)(2S_2 + 1)}{1(2S_4 + 1)} \times \frac{k_{\alpha}^2}{k_{\gamma}^2} \times \sigma(\alpha, \overrightarrow{\gamma_i}) \end{aligned}$ $\text{but: } \sigma(\alpha, \overrightarrow{\gamma_i}) = \frac{1}{2}\sigma(\alpha, \gamma) \\ & = \frac{(2S_1 + 1)(2S_2 + 1)}{2(2S_4 + 1)} \times \frac{k_{\alpha}^2}{k_{\gamma}^2} \times \sigma(\alpha, \gamma) \\ & = \frac{1}{2} \times (80 - 160) \times \sigma(\alpha, \gamma) \end{aligned}$

$$= (40 - 80) \times \sigma(\alpha, \gamma)$$

Luminosity = 1 nb⁻¹ per Day $I_{\gamma} = 4.0 \times 10^7$ /sec (phase 1), $\Delta E = 2\%$, Collimator = 1/2" $\sigma_{\alpha\gamma}(1.3MeV) = 0.6 \ nb, \ \sigma_{\gamma\alpha}(8.5MeV) = 30 \ nb \rightarrow 30CPD$





One tantalizing new problem that was posed by the new gamma-ray data [2] is the disagreement between E1/E2 mixing phases (ϕ_{12}) extracted from the measured gamma-ray angular distributions and the mixing phase predicted by theory:

 $\phi_{12} = \delta_2 - \delta_1 + \arctan(\eta/2)$

where δ_1 and δ_2 are the p and d wave elastic phase shifts and η is the Sommerfeld parameter. In Ref. [2] the disagreement is considered as a simple disagreement between data and the prediction of R-Matrix theory. But in fact the above relationship is rooted in the Watson theorem and unitarity. This disagreement is observed already at high energies, on the broad 1⁻ resonance located at 9.58 MeV where the capture cross-section is large, and it must be resolved.























L. Weissman et al.; JINST 1(2006)05002







Quantum Efficiency of CHORUS BV55 Optical Chain/ PTB. March 2006.

Contrast Transfer Function Mask, PTB, June 2006.





Fig. 56: ROI full vertical binning of the 3 mm mask (a = white value; b = black value; c =pixel contents)





Benjamin Bromberger - PTB (University of Braunschweig)

X

TOPAC2



Mohammad W. Ahmad - TUNL

Triggered Readout of OTPC System



BUSY OR



The Entire System

















Lens Fiducial Circle (D ~ 15cm) at 85 cm Distance

Data Profile

























Channel

<u>Horizontal Track</u>



Slightly Tilted Track

lek "ľ	1 🛑 Acq C	Complete M Pos: 0.000s	CH2
		. (* 4	Coupling
			BW Limit 200MHz
			Volts/Div
			Probe 1008
			Invert IIII
H1 20.0mV		M 1.00.05 CH2 /	2. 4 8V

Very Tilted Track

CH1 500mV		 	-								2 48V
		 	Ξ.,								Invert
		 									Probe 10X
											Volts/D Coarse
				 							Off 200MH
1		,			aj). Int					11 - 101 11 - 101	
Tek		 Acq	Con	nple	ite	M	Pos	: 1.3	20,	JS	CH2

Tracks From ¹⁴⁸Gd (3.181 MeV)









Laser power, low emittance, low bandwidth puts T-REX at the forefront of high brightness



We estimate T-REX peak brilliance at 1 MeV exceeds synchrotrons by 15 orders of magnitude

University of Connecticut Laboratory for Nuclear Science at Avery Point

<u>August, 2007:</u>	O-TPC moves to TUNL (We need a clean room)
<u>September, 2007:</u>	O-TPC calibrated at TUNL DAQ system tested
<u>Fall, 2007:</u>	One shift test <u>Phase 1:</u> First week of data Break Second week of data
	Supplementary Request (DOE ~\$40K)
<u>January 21, 2008:</u>	End of my sabbatical leave!!!
<u>Summer 2008:</u>	Please Consider an upgrade of HI y S ala T-REX (LLNL)